

## Spatial Weighting and Iterative Projection Methods for EOFs

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### ABSTRACT

Often there is a need to consider spatial weighting in methods for finding spatial patterns in climate data. The focus of this paper is on techniques that maximize variance, such as empirical orthogonal functions (EOFs). A weighting matrix is introduced into a generalized framework for dealing with spatial weighting. One basic principle in the design of the weighting matrix is that the resulting spatial patterns are independent of the grid used to represent the data. A weighting matrix can also be used for other purposes, such as to compensate for the neglect of unrepresented subgrid-scale variance or, in the form of a prewhitening filter, to maximize the signal-to-noise ratio of EOFs. The new methodology is applicable to other types of climate pattern analysis, such as extended EOF analysis and maximum covariance analysis. The increasing availability of large datasets of three-dimensional gridded variables (e.g., reanalysis products and model output) raises special issues for data-reduction methods such as EOFs. Fast, memory-efficient methods are required in order to extract leading EOFs from such large datasets. This study proposes one such approach based on a simple iteration of successive projections of the data onto time series and spatial maps. It is also demonstrated that spatial weighting can be combined with the iterative methods. Throughout the paper, multivariate statistics notation is used, simplifying implementation as matrix commands in high-level computing languages.

### 1. Introduction

The analysis of spatial patterns in geophysical data is performed with a wide variety of techniques, including methods based on empirical orthogonal functions (EOFs; e.g., von Storch and Zwiers 1999; Jolliffe 2002; Bretherton 2003; Hannachi et al. 2006, 2007; Van den Dool 2007), empirical normal modes (Brunet 1994; Brunet and Vautard 1996), methods that find coupled patterns between datasets (e.g., maximum covariance analysis, canonical correlation analysis, combined principal component analysis; see von Storch and Zwiers 1999), methods that find similar patterns, such as cluster analysis (e.g., Cheng and Wallace 1993; Wallace et al. 1993a) and self-organized maps (Hewitson and Crane 2002; Reusch et al. 2007). With any of these methods, the use of spatial weighting may be an important issue for several reasons.

One reason to use spatial weighting is that spatial patterns should be invariant to how one chooses the grid locations, since one is aiming to find properties of the continuous spatial field (e.g., “intrinsic EOFs”; North et al. 1982; Stephenson 1997). Spatial weighting can be used to compensate for unequal distribution of grid points. Weighting may be desirable for other purposes, such as to emphasize (or mask) certain spatial regions, to account for variations in error covariances, to calculate patterns with more than one variable measured across the spatial domain (e.g., extended EOF analysis), or to equalize the variance at every grid point (e.g., EOFs based on the correlation matrix instead of the covariance matrix). Weighting may also be used to compensate for small-scale variance that is not represented by the gridded data matrix. It may also be necessary to apply weighting to find EOFs of quantities such as zonally averaged angular momentum, derived from the zonal-mean zonal wind field (e.g., Baldwin and Tung 1994). Weighting may also take the form of a prewhitening filter (Allen and Smith 1997; Venzke et al. 1999; Chang et al. 2000) in which the data matrix is premul-

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tiplied by a filter (weighting matrix) that is determined by the noise covariance matrix.

In this paper we adopt a general approach to spatial weighting of geophysical fields. We use EOF analysis as an example to introduce a generalized weighting matrix. We show that it is possible to include the weighting matrix in the EOF calculation, without premultiplying the data by the square roots of the weights, as was done by North et al. (1982). Our results are applicable to a wide variety of similar techniques—in general any techniques that partition variance.

It is often desirable to perform EOF analysis on large, possibly 3D, spatiotemporal datasets (e.g., Hawkins and Sutton 2007). Standard EOF techniques can become impractical for large data matrices. Efficient techniques such as the power method can be used to find the leading EOFs (Golub and Van Loan 1983; Jolliffe 2002 for a review; Van den Dool 2007) but typically involve the explicit calculation of the whole covariance matrix, which can make even these techniques impractical. It becomes necessary to have EOF methods that can rapidly find the leading EOFs for such datasets while allowing for arbitrary spatial weighting.

This paper addresses these issues by 1) proposing a simple iterative scheme suitable for finding EOFs in large datasets, and 2) mathematically formulating the EOF problem for arbitrary weighting schemes. The iterative method is applicable to any of the EOF-based techniques and uses the data matrix directly. We demonstrate that the convergence properties of the iterative method are the same as those of the power method (Golub and Van Loan 1983).

The outline of the paper is as follows. Section 2 introduces the nomenclature and fundamentals of EOFs in standard matrix notation. Section 3 discusses spatial weighting in EOF analysis; in section 3a we introduce a weighting matrix, and in section 3b we discuss how to deal with subgrid-scale variance. Section 3c presents a derivation of a generalized weighting metric, which can be easily transformed away by a change of variables involving multiplying the data by the square root of the weighting metric. Section 4a discusses in general how to “project” data onto time series and spatial patterns. We then show, in section 4b, how an iterative projection method can be combined with spatial weighting to efficiently calculate the leading EOF. In section 5 we provide examples of including weighting matrices in EOF analysis.

## 2. Nomenclature and approach to EOF analysis

We begin by considering a field, such as pressure or temperature that is a continuous function of space and time. We assume that the continuous field is approxi-

mated by sampling at regular time intervals onto some spatial grid, which may or may not be regular. The data are organized in an  $(n \times p)$  matrix,  $\mathbf{X} = [\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_n]^T$ , containing  $n$  observations of a variable  $\mathbf{x}_t$  defined at  $p$  spatial points. The  $p$  spatial points could be embedded in one, two, or three spatial dimensions. The  $n$  observations could be of the same field made at different times or  $n$  observations of different images (e.g., images of human faces; Craw and Cameron 1992). For the purposes of this paper we will assume, without loss of generality, that we are dealing with  $p$  spatial observations at  $n$  regularly spaced times.

We also assume the variables in  $\mathbf{X}$  (the time series, or columns) are anomalies that have been obtained by centering the original variables by removing the respective time means at each location. The means of each time series (i.e., the columns of  $\mathbf{X}$ ) are therefore identically zero. Because the variables are centered, the  $(p \times p)$  sample covariance matrix is given by the matrix product  $\mathbf{S} = \mathbf{X}^T \mathbf{X} / (n - 1)$ . Here it is conventional to use  $n - 1$  (instead of  $n$ ) when the covariance is computed from a sample. The *total variance* is the sum of the variances at each of the grid points (the diagonal elements of  $\mathbf{S}$ ) and so is given by the trace of  $\mathbf{S}$ ,  $\text{Tr}(\mathbf{S})$ .

EOF analysis decomposes a data matrix into a series of data-based orthogonal functions, which are defined so that a minimal number of EOFs are needed to reconstruct the variation within the original data matrix. EOF analysis can be understood in terms of spatial patterns (EOFs) and their associated time series (often called principal components). The data matrix can be written as a sum of the products of the EOFs  $[\mathbf{e}_1, \mathbf{e}_2, \dots]$  and their associated time series  $[\mathbf{y}_1, \mathbf{y}_2, \dots]$ :

$$\mathbf{X} = \sum_{i=1}^r \mathbf{y}_i \mathbf{e}_i^T, \quad (2.1)$$

where  $r$  is the rank of  $\mathbf{X}$ , which is never greater than the minimum of  $n$  and  $p$ . The EOFs  $[\mathbf{e}_1, \mathbf{e}_2, \dots]$  are chosen so as to be orthonormal ( $\mathbf{e}_i^T \mathbf{e}_j = 1$  for  $i = j$  and 0 otherwise) and successively maximize the variance in the corresponding time series  $[\mathbf{y}_1, \mathbf{y}_2, \dots]$ . Here  $\mathbf{e}_1^T \mathbf{S} \mathbf{e}_1$  is maximal and each successive EOF component accounts for the greatest possible fraction of the remaining variance. The EOF time series can be shown to be uncorrelated with one another, but this does not necessarily mean that they are independent. For example, a pair of EOFs may describe a propagating pattern in a data matrix [e.g., the quasi-biennial oscillation (QBO), as shown by Wallace et al. (1993b)].

## 3. Spatial weighting in EOF analysis

One common reason to weight in EOF analysis is to compensate for grid spacing in the data matrix, for ex-

ample to compensate for converging meridians on a latitude–longitude grid. But there are several other reasons to consider weights in EOF analysis. It may be necessary to create a metric of some integral quantity of interest, such as forming angular momentum from zonal wind data or weighting atmospheric data in the vertical by inverse pressure or by height. It may be desirable to apply different weights to each variable when several variables are included in the analysis. In some cases one may wish to emphasize or de-emphasize certain geographic regions or data levels. Weighting may be used to correct for spatially inhomogeneous data variances or error covariances. EOF analysis is often done using correlations rather than covariances, which corresponds to weighting anomalies by the reciprocal of their respective standard deviations.

As discussed in North et al. (1982), there is typically a finite number of sometimes irregularly spaced grid points in a domain. We wish to find what North et al. termed the “intrinsic EOFs,” which are independent of the data grid. The concept of intrinsic patterns, independent of the grid, can be applied more broadly to any type of climate pattern analysis. To do this we must find an appropriate weighting that is invariant to the choice of grid spacing.

We will use EOF analysis as an example, but the resulting weighting methods are broadly applicable to all variants of EOF methods (e.g., empirical orthogonal teleconnections; Van den Dool et al. 2000). This weighting method is also applicable to maximum covariance analysis (MCA) because MCA with identical input fields finds the EOFs. For any of these methods, the weighting technique must produce the same patterns and time series as would result from regridding the data onto an equal-area grid.

Despite all these reasons, the subject of how to weight EOFs has received little formal attention in the published scientific literature. This section will attempt to address this deficit by showing how EOF analysis can be generalized to include arbitrary weighting metrics.

#### a. Weighting metric

The generalization of the EOF problem to continuous (Karhunen–Loève) functions was discussed by North et al. (1982). They reasoned that EOFs derived from an approximation to the continuous field, on a finite set of grid points, should approximate those based on the continuous functions, and should therefore be independent of the particular grid chosen. Our approach is similar in that we begin by considering a continuous anomaly field.

It is important to note that EOF methods partition the total variance of the data matrix, which is a sum of

squared anomalies from the mean. For the purpose of introducing a weighting matrix, it is sufficient to consider a continuous anomaly field  $\mathbf{x}$  at a single time, such as sea level pressure over the Northern Hemisphere. The area integral of  $\mathbf{x}^2$  would correspond to a summation of squared gridpoint values. If  $\mathbf{x}$  is approximated on a series of  $p$  grid points, then the area integral of the squared continuous anomaly field  $\mathbf{x}^2$  can be approximated by the quadratic sum of anomaly values at the grid points  $i$ . If the grid is uniform, and therefore weighting is not needed,

$$\int_A \mathbf{x}^2 da \approx \sum_{i=1}^p \mathbf{x}_i^2 = \mathbf{x}^T \mathbf{x}. \quad (3.1)$$

To accommodate the most general types of weighting, we introduce the  $(p \times p)$  weighting matrix  $\mathbf{W}$ :

$$\int_A \mathbf{x}^2 da \approx \sum_{i=1}^p \sum_{j=1}^p x_i W_{ij} x_j = \mathbf{x}^T \mathbf{W} \mathbf{x}. \quad (3.2)$$

We define the weighting matrix in a general way so that it can be used to area weight the summation, or to weight or filter EOFs in other circumstances (described below). It is a generalization of the area-weighting concept discussed in North et al. (1982), who defined a  $p$ -vector of weights proportional to the gridbox area.

The weighting matrix can always be written symmetrically  $W_{ij} = W_{ji}$  without any loss of generality. Since the integral of a squared quantity is nonnegative, it is desirable that  $\mathbf{x}^T \mathbf{W} \mathbf{x} \geq 0$  and so the weighting matrix  $\mathbf{W}$  then defines a metric  $\|\mathbf{x}\|_{\mathbf{W}}^2 = \mathbf{x}^T \mathbf{W} \mathbf{x}$  in the  $p$ -dimensional space of gridpoint variables (Strang 2003).

It is common practice to assume that the area integral can be approximated by the simple sum of squared anomalies at each grid point—in other words, the midpoint rule for numerical integration. In the simplest case of uniform grid spacing,  $\mathbf{W}$  is the identity matrix, with elements  $W_{ij} = \delta_{ij}$ , where  $\delta_{ij} = 1$  only if  $i = j$  and is 0 otherwise. For nonuniform grid spacing, it is common to modify this by assuming a diagonal weighting matrix with  $W_{ij} = a_i \delta_{ij}$ , where  $a_i$  is the area associated with the  $i$ th grid point. This corresponds to the equal-area weighting discussed by North et al. (1982).

Defining  $\mathbf{W}$  as a  $(p \times p)$  matrix allows for the inclusion of off-diagonal elements. The  $\mathbf{W}$  has the same dimensions as the covariance matrix and so allows for the weighting to vary not just by grid position but to vary across the covariance matrix. In some cases more accurate estimates of the area integral can be obtained by using sophisticated numerical integration together with careful spatial smoothing that takes account of the spatial correlations in  $\mathbf{x}^2$  (Kagan et al. 1997). The result of such approaches can still be written in the form  $\mathbf{x}^T \mathbf{W} \mathbf{x}$ ,

where the effect of smoothing and spatial interpolation is to induce off-diagonal elements in  $\mathbf{W}$ . The effects of spatially varying errors in  $\mathbf{x}^2$  can also be taken into account by including the inverse of an error covariance matrix in  $\mathbf{W}$ , as is done in optimal interpolation and kriging (Kagan et al. 1997).

It is important to note that any metric  $\mathbf{W}$  can always be transformed away to Euclidean (uniform weighting) by pretransforming the gridpoint variables to new gridpoint variables:

$$\mathbf{x} \rightarrow \mathbf{x}' = \mathbf{W}^{1/2}\mathbf{x}. \quad (3.3)$$

For diagonal metrics, such as the area-weighting metric,  $\mathbf{W}^{1/2}$  is simply the square root of the diagonal elements of  $\mathbf{W}$ . For nondiagonal metrics,  $\mathbf{W}^{1/2}$  can be calculated by using numerical methods such as the Cholesky decomposition or the Denman–Beavers iterative method (Higham 2008).

The generalized norm can then be written as a Euclidean norm in the new variables:

$$\mathbf{x}^T \mathbf{W} \mathbf{x} = (\mathbf{W}^{1/2}\mathbf{x})^T (\mathbf{W}^{1/2}\mathbf{x}) = \mathbf{x}'^T \mathbf{x}'. \quad (3.4)$$

This is the reason for premultiplying anomalies by the *square root* of the gridbox area, as suggested by North et al. (1982). Alternatively, one can find transformations of gridpoint variables that can remove the metric, for example, by interpolating  $x$  to a grid with equal areas (or volumes).

Although the metric  $\mathbf{W}$  was introduced in order to compensate for unequal grid spacing, the concept can be used in a variety of ways to weight or filter data in EOF analysis. One example is to prewhiten the data in signal-to-noise maximizing EOF analysis (Allen and Smith 1997; Venzke et al. 1999; Chang et al. 2000). In this case the data matrix is premultiplied by a filter (in place of  $\mathbf{W}^{1/2}$ ) that is completely determined by the noise covariance matrix. If, for example, the analysis was on a nonuniform grid, the filter matrix and  $\mathbf{W}^{1/2}$  could be combined by multiplying the data matrix by both and then performing the EOF analysis.

Another use of the weighting metric is to emphasize (or de-emphasize) regions of the data matrix. This may be desirable because of signal-to-noise ratio, to produce uniform variance across the grid, or simply to eliminate (mask) a region from the analysis. In all these cases, once  $\mathbf{W}$  is defined, it is used in the same way in the EOF analysis.

#### *b. Choice of metric to account for subgrid-scale variation*

A continuous squared anomaly field  $\mathbf{x}^2$  can be approximated on grid points, but doing so acts as a spatial

smoother. This reduction in variance may be small if the decorrelation distance is large, but the reduction can be significant if the field has large small-scale variance. The size of the grid spacing therefore affects the total variance as well as the EOFs. This suggests that EOFs should be calculated on the highest grid resolution available, if one wishes to include the effects of small-scale variance. There may be instances in which it may be desirable to calculate EOFs of large-scale variability [in analogy with EOFs of “low-frequency variability,” as in Thompson and Wallace (1998)]. For many datasets the decorrelation distance may be highly variable over different regions of the grid. If EOFs are calculated on a variable grid (e.g., a latitude–longitude grid), there are two reasons that weighting may be necessary: first to compensate for the gridbox area and second to compensate for unrepresented variance at scales smaller than the grid.

Compensation for unrepresented variance can be subtle, in that the coarseness of the grid affects how much of the variance of the continuous field can be represented by the grid. Some of the variance of the continuous field is not represented by the gridpoint values. In effect, there is variation of the continuous field within each grid box, so that the squares of the data values are underrepresented if the data vary within a grid box.

A continuous anomaly field  $x$  could vary within a grid box because of the presence of spatial trends across grid points and subgrid-scale random variations in  $\mathbf{x}^2$  within grid boxes. However, we can only estimate the area integral of  $\mathbf{x}^2$  using gridpoint values  $\mathbf{x}_i$  that are considered to be area averages of the field across the grid box. By Reynolds averaging (Monin and Yaglom 1971), it can be seen that the area integral of  $\mathbf{x}^2$  across a grid box exceeds the area weighting commonly used in climate studies. If we integrate  $\mathbf{x}^2$  over a single grid box,

$$\int_A \mathbf{x}^2 da = \mathbf{x}_i^2 a_i + \int_A (\mathbf{x} - \mathbf{x}_i)^2 da. \quad (3.5)$$

The first term on the right of (3.5) represents the area weighting discussed in section 3a. The second term represents variance at scales smaller than the grid box. If the grid is fine enough so that the gridded values account for nearly all the variance, the second term is small and area weighting  $a_i$  would be a good approximation. If the grid is too coarse, the first term does not adequately account for the total variance of the field.

To illustrate this effect, we consider a single Northern Hemisphere grid of the 40-yr European Centre for Medium-Range Weather Forecasts (ECMWF) Re-Analysis (ERA-40) geopotential (Uppala et al. 2005) at

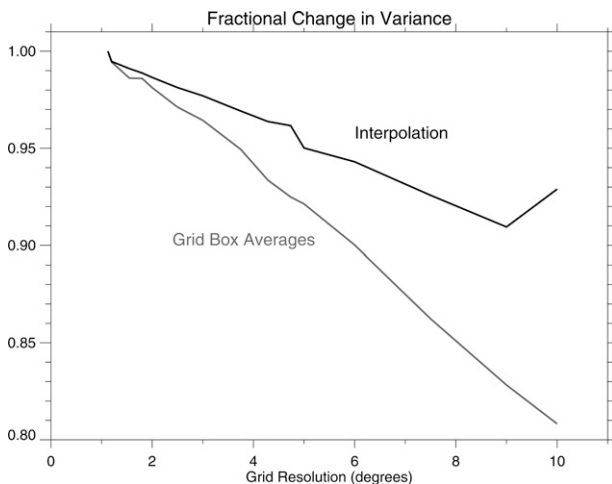


FIG. 1. Fractional change in Northern Hemisphere average variance of 1000-hPa geopotential on 1 Jan 2001. The full grid resolution is 1.125°. The black curve shows the total variance at coarser grid resolutions obtained by bilinear interpolation. The gray curve is obtained by averaging over grid boxes.

1000 hPa for 1 January 2001. The original resolution is a 1.125° latitude–longitude grid. Figure 1 examines the ratio of  $\sum_i \mathbf{x}_i^2 a_i$  at lower grid resolutions to  $\sum_i \mathbf{x}_i^2 a_i$  using the original 1.125° grid. We use two methods for calculating reduced-resolution grid values: 1) interpolation and 2) spatial averages over each grid box. Interpolated values (solid black curve) in Fig. 1 illustrate that, as the grid resolution is decreased, the mean grid variance is decreased to ~95% at 5°–6° and ~90% at 10°. Gridbox averaging (gray curve) tends to reduce local minima and maxima in the grid (compared to interpolation) so the falloff is somewhat steeper. Figure 1 illustrates that ~10%–20% of the ERA-40 1000-hPa geopotential variance is at grid resolutions smaller than 10°.

In principle, the most extreme case of small-scale variance would be a random anomaly field with variance  $\sigma^2$  and no spatial correlation between grid points. The area integral would be equal to  $\sigma^2 a_i$  but would have a squared gridbox average that scales as  $\mathbf{x}_i^2 = \sigma^2/a_i$  (the variance of a mean). In this special case, it would be possible to compensate for the subgrid-scale variance by increasing the weighting factor to  $a_i^2$  rather than  $a_i$  (Folland 1988). Thus, the correct weighting factor for any field would lie between  $a_i$  and  $a_i^2$ , but for most data fields (e.g., the geopotential used in Fig. 1), the factor would be close to  $a_i$ .

To further illustrate this effect, consider a uniformly gridded data matrix in which the  $\mathbf{x}_i$  are uncorrelated with each other and have unit variance. If half of the grid is resampled with grid boxes twice as large, each new gridpoint value would represent the average of two original values ( $a_i = 2$ ). The variance of the average of

two such uncorrelated times series is reduced from 1 to 0.5. Thus, a weighting factor of 2 would be necessary to compensate for the gridbox area, and another factor of 2 would be needed to compensate for the reduced variance in the resampled half of the grid, giving  $a_i^2 = 4$  as the correct weighting, as found by Folland (1988).

Defining a weighting matrix that compensates for unrepresented variance within grid boxes is possible, at least in principle, and would depend on the decorrelation scale and power spectrum of the data at length scales smaller than the grid boxes. For grid spacing smaller than the decorrelation scale of the field (e.g., sea level pressure on a grid much smaller than the synoptic scale), the correct metric would be  $W_{ij} = a_i \delta_{ij}$  (area weighting), whereas, for grid boxes containing less coherent variations, the limit of  $W_{ij} = a_i^2 \delta_{ij}$  (weighting by the square of the gridbox area) could be approached. The correct choice of metric would ensure that the results of the EOF analysis would not be overly dependent on the specific choice of grid.

### c. Generalized EOF analysis

The time average of the numerical approximation to the area integral can be written in terms of the sample covariance as follows:

$$\overline{\int \mathbf{x}^2 da} = \int \overline{\mathbf{x}^2} da \approx \frac{1}{n} \sum_{t=1}^n \mathbf{x}_t^T \mathbf{W} \mathbf{x}_t = \text{Tr}(\mathbf{W} \mathbf{S}), \quad (3.6)$$

which defines the generalized total variance. The EOF problem becomes one of finding spatial patterns that maximize  $\mathbf{e}_i^T \mathbf{W} \mathbf{S} \mathbf{e}_i$  subject to the constraint  $\mathbf{e}_i^T \mathbf{W} \mathbf{e}_j = 1$  for  $i = j$  and 0 otherwise. By transforming to  $\mathbf{e}'_i = \mathbf{W}^{1/2} \mathbf{e}_i$ , it can be seen that the EOF solutions are the eigenvectors of the generalized eigenvector equation  $\mathbf{S} \mathbf{W} \mathbf{e}_i = \lambda_i \mathbf{e}_i$ .

Several different approaches can be used to find the generalized EOFs:

- 1) Solve the generalized eigenvector equation ( $\mathbf{S} \mathbf{W} \mathbf{e}_i = \lambda_i \mathbf{e}_i$ ) directly [e.g., using generalized solvers such as the eig function in matrix laboratory (MATLAB)].
- 2) Transform out the metric by pretransforming the original variables, find unweighted EOFs, and then back transform the resulting EOFs. Pretransforming maps the data matrix to  $\mathbf{X}' = \mathbf{X} \mathbf{W}^{1/2}$  and the covariance matrix to  $\mathbf{S}' = (\mathbf{W}^{1/2})^T \mathbf{S} \mathbf{W}^{1/2}$ . It is necessary to perform the inverse transformation  $\mathbf{e} = \mathbf{W}^{-1/2} \mathbf{e}'$  to obtain the generalized EOFs from the EOFs of the transformed variables.
- 3) Use iterative projection methods that include  $\mathbf{W}$  explicitly (section 4).

#### 4. Projection methods for finding EOFs

##### a. Projections and EOFs

For any  $n$ -vector time series of centered anomalies  $\mathbf{y}$ , an associated  $p$ -vector spatial pattern  $\mathbf{e}$ , can be obtained by “projecting the data onto the time series,” as follows:

$$\mathbf{e} = \frac{\mathbf{X}^T \mathbf{y}}{\mathbf{y}^T \mathbf{y}}, \quad (4.1)$$

where  $\mathbf{y}^T \mathbf{y} = (n-1)s_y^2$  and  $s_y$  is the sample standard deviation of the time series. For example,  $\mathbf{X}$  could be monthly-mean Northern Hemisphere sea level pressure data, and  $\mathbf{y}$  could be an index of the North Atlantic Oscillation (NAO). Equation (4.1) gives the spatial pattern of the NAO. Note that when  $\mathbf{y} = \mathbf{x}_i$ , the time series at grid point  $i$ ,  $e_j = s_{ji}/s_{ii}$  is the covariance map at grid point  $i$ .

Similarly, a time series  $\mathbf{y}$  can be obtained from a spatial pattern  $\mathbf{e}$  by “projecting the data onto the spatial pattern”:

$$\mathbf{y} = \frac{\mathbf{X}\mathbf{e}}{\mathbf{e}^T \mathbf{e}}. \quad (4.2)$$

The metric  $\mathbf{W}$  does not affect the projection Eq. (4.1) for obtaining a spatial pattern from a time series, but Eq. (4.2) for obtaining an EOF time series from an EOF requires weighting of both  $\mathbf{X}$  and  $\mathbf{e}$  by  $\mathbf{W}^{1/2}$ :

$$\mathbf{y} = \frac{\mathbf{X}\mathbf{W}\mathbf{e}}{\mathbf{e}^T \mathbf{W}\mathbf{e}}. \quad (4.3)$$

For example, a daily NAO index can be obtained from the NAO spatial pattern and daily sea level pressure anomalies (e.g., Baldwin et al. 2003) using (4.3). Compensating for a latitude–longitude grid would require the diagonal elements of  $\mathbf{W}$  to be  $\cos \theta$ . Note that, if the data matrix is pretransformed by multiplying by  $\mathbf{W}^{1/2}$ , the resulting EOFs will also be transformed. To obtain the EOFs for the original data matrix, either use (4.1) or divide by  $\mathbf{W}^{1/2}$ .

Normalization between the spatial patterns and time series is arbitrary, as long as the product  $\mathbf{y}\mathbf{e}^T$  is preserved. If the spatial patterns are treated as dimensionless weights, then  $\mathbf{e}^T \mathbf{W}\mathbf{e}/\text{Tr}(\mathbf{W}) = 1$ . If the EOF time series have unit variance, then  $\mathbf{y}^T \mathbf{y}/(n-1) = 1$ . If the EOFs are normalized to have unit normalization, then  $\mathbf{e}^T \mathbf{W}\mathbf{e}/[(p-1)\text{Tr}(\mathbf{W})] = 1$ .

##### b. A simple iterative projection method for calculating EOFs

Several numerical matrix algebra techniques, such as eigenvector analysis and singular value decomposition

(SVD) can be used to obtain all the EOFs simultaneously (Jolliffe 2002). For large data grids, these methods can rapidly become computationally impractical. Although the space and time dimensions in the data matrix can be swapped (von Storch and Zwiers 1999) even calculating the covariance matrix can be problematic. One alternative is to use SVD directly on the data matrix, rather than the covariance matrix, but the SVD method is still computationally expensive for large data matrices.

Rather than calculating all the EOFs, the first few EOFs and EOF time series can be calculated using iterative techniques (e.g., Holmström 1963; Clint and Jennings 1970; Jongman et al. 1995; Legendre and Legendre 1998). One such technique, called the “power method” (Jolliffe 2002), begins with an initial estimate of the leading EOF, with the only requirement being that the initial estimate has at least some projection onto the leading EOF. The estimated EOF is repeatedly multiplied by the covariance matrix  $\mathbf{X}^T \mathbf{X}$  and normalized during each iteration:

$$\mathbf{e}_{k+1} = \frac{\mathbf{X}^T \mathbf{X}\mathbf{e}_k}{\|\mathbf{X}^T \mathbf{X}\mathbf{e}_k\|}. \quad (4.4)$$

The successive estimates converge to the leading EOF as long as the initial estimate has at least minimal projection onto the leading EOF (Golub and Van Loan 1983, p. 209). The rate of convergence depends on the ratio of the two leading eigenvectors. Thus, an arbitrary initial guess at the EOF is sufficient in practice to initialize the procedure. The EOF time series is never used explicitly, but it would be obtained using (4.2). Once convergence is achieved, the projection of the leading EOF ( $\mathbf{y}\mathbf{e}^T$ ) is subtracted from  $\mathbf{X}$ , and the procedure is repeated to find the desired number of EOFs.

Here we describe another iterative method of finding the leading EOF that involves successive application of (4.1) and (4.3). We will show that this method is an extension of the power method. The main advantages are that this iterative procedure does not require the calculation of the covariance matrix  $\mathbf{X}\mathbf{X}^T$ , and it accounts for the weighting matrix  $\mathbf{W}$ . This method performs the multiplication by  $\mathbf{X}\mathbf{X}^T$  in two steps and so retains the convergence properties of the power method. The method consists of a simple iterative procedure using the projections (4.1) and (4.3), which alternate between the EOF and its time series. By successively multiplying by  $\mathbf{X}$  and  $\mathbf{X}^T$ , we calculate both the EOF and its time series (which was implicit in the power method). Another method involving iteration between a spatial pattern and a time series was proposed by Clint and Jennings (1970). Van den Dool et al.

(2000) used a similar approach, without weighting, to find the leading EOF beginning from the leading empirical orthogonal teleconnection (EOT) pattern. Iteration between a time series and spatial pattern to calculate the leading EOF was discovered independently by G. Hegerl (2008, personal communication).

The method begins with an initial random guess time series whose only requirement is that (as with the power method) it must have a nonzero projection onto the leading EOF time series. Let  $\mathbf{y}_k$  denote the  $k$ th iteration of the first EOF time series. The data matrix is then projected onto this time series to obtain a spatial pattern, followed by projecting the data matrix onto this spatial pattern to get a new time series. The process is then iterated until the squared error between successive time series falls below a preset tolerance, which can be small but must be greater than the machine precision:

$$\mathbf{e}_{k+1} = \frac{\mathbf{X}^T \mathbf{y}_k}{\mathbf{y}_k^T \mathbf{y}_k} \quad \text{and} \quad (4.5)$$

$$\mathbf{y}_{k+1} = \frac{\mathbf{X} \mathbf{W} \mathbf{e}_{k+1}}{\mathbf{e}_{k+1}^T \mathbf{W} \mathbf{e}_{k+1}}. \quad (4.6)$$

Iteration continues until convergence, defined by  $\|\mathbf{y}_{k+1} - \mathbf{y}_k\|^2 < \varepsilon$ , where  $\varepsilon$  is a preset tolerance. The denominators in (4.5) and (4.6) are normalization factors. Substituting (4.6) into (4.5) demonstrates that this method is equivalent to the power method. Once convergence is achieved, the EOF multiplied by its time series is subtracted from the data matrix  $\mathbf{X} = \mathbf{X} - \mathbf{y} \mathbf{e}^T$ , as with the power method, and the process is repeated to find the next EOF. There is no advantage in transposing the data matrix (switching time and space). After convergence the EOF and its time series can be renormalized in any way that preserves the product  $\mathbf{y} \mathbf{e}^T$ .

The algorithm finds one EOF at a time, so it is appropriate to use when only a small number of EOFs are desired. If the eigenvalues of the leading modes are not well separated, convergence will be slow (Golub and Van Loan 1983, p. 209). Since this iterative approach does not explicitly compute the covariance matrix, it can therefore be used when  $\mathbf{X}$  is a very large matrix. The algorithm does not require significant memory beyond that needed for the data matrix, and the algorithm would work even if the data matrix exceeds computer memory. For example, we have used as a data matrix of 50 yr of daily Northern Hemisphere ERA-40 geopotential (dimensions 18 262 by 25 920). The calculation of the leading EOFs was routine using the iterative algorithm.

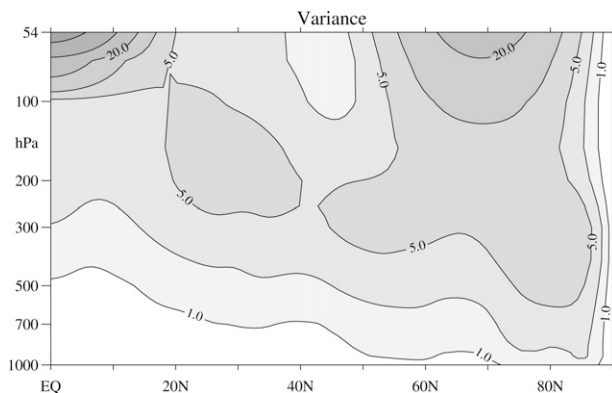


FIG. 2. Variance of monthly-mean, ERA-40 Northern Hemisphere deseasoned, detrended zonal-mean zonal wind, 1958–2001. The contour intervals are 1, 2.5, 5, 10, 20, 40, 80  $\text{m}^2 \text{s}^{-2}$ . The highest level is 54.6 hPa (one of the ERA-40 assimilation model levels). The data levels correspond to the ticks on the right axis.

## 5. Example

The purpose of this section is to illustrate how to create a weighting matrix and how different choices for a weighting matrix affect the resulting EOF. It should be emphasized that resulting EOFs may or may not represent physical modes (e.g., Brunet and Vautard 1996; Ambaum et al. 2001; Jolliffe 2003). The strict orthogonality condition on EOFs is at the same time its strength and its weakness. For example, the North Atlantic Oscillation may be a more physical mode than the similar EOF of the northern annular mode (Ambaum et al. 2001). We consider the leading EOF of Northern Hemisphere monthly-mean zonal-mean zonal wind from 1000 to 54 hPa. The data are ERA-40 (Uppala et al. 2005) from 1958 to 2001 on a  $1.125^\circ$  grid in latitude and unequally spaced pressure levels. The dimensions are  $n = 528$  months and with 81 latitudes and 15 pressure levels,  $p = 1215$ . To remove the annual cycle and the long-term trend, we subtracted the average value for each calendar month and then regressed out the long-term linear trend at each grid point. Figure 2 illustrates the variance  $\bar{\mathbf{x}}^2$  of the resulting anomalies. The variance is largest in the stratosphere, both near the equator, where the quasi-biennial oscillation dominates (Baldwin et al. 2001), and at high latitudes where month-to-month wintertime changes in zonal wind are large.

This small dataset allows us to illustrate how to create three different weighting matrices that varies in both latitude and pressure. We choose  $\mathbf{W}$  to be a diagonal matrix (one weighting value per grid point), so the diagonal elements of  $\mathbf{W}$  map to the grid points in the latitude–height plane. To weight in both latitude and pressure, we specify both a latitudinal weighting

function and a vertical weighting function. Then  $\mathbf{W}$  is formed by multiplying these functions. Latitudinal weighting of  $\cos\theta$  compensates for the converging meridians at higher latitudes. In the vertical, we choose two different weightings. The first is proportional to the difference in pressure (or mass) between the bottom and top of the layer represented by each level. In the first case,  $W_{ij} = \cos\theta_i \Delta p_i \delta_{ij}$ . The second case is proportional to the difference in log pressure (approximately physical height):  $W_{ij} = (\cos\theta_i \Delta p_i \delta_{ij})/p_i$ . These weightings are illustrated in Figs. 3a,b. As a third example, we weight the zonal-mean wind to form a quantity proportional to axial relative angular momentum by pressure weighting in the vertical and  $\cos^2\theta$  weighting in latitude (Fig. 3c). The global axial relative angular momentum is given by (von Storch 2000)

$$M_r = \frac{\pi r^3}{4} \int_{-\pi/2}^{\pi/2} \int_0^{P_s} (\bar{u} \cos^2\theta) dp d\theta, \quad (5.1)$$

where  $r$  is earth's radius and  $P_s$  is surface pressure. Relative angular momentum is therefore proportional to  $\bar{u} \cos^2\theta$  and  $W_{ij} = \cos^2\theta_i \Delta p_i \delta_{ij}$ . This weighting is illustrated in Fig. 3c.

The magnitude of  $\mathbf{W}$  does not affect the EOFs. The weighting matrix can be multiplied by an arbitrary scaling value without affecting the EOFs. The leading EOFs, corresponding to the weightings in Fig. 3, are shown in Fig. 4. The EOFs were obtained by iterative projections and include  $\mathbf{W}$  explicitly in the projection used to find the time series [Eqs. (4.5) and (4.6)]. The same result could have been obtained by premultiplying by  $\mathbf{W}^{1/2}$ . However, premultiplication by  $\mathbf{W}^{1/2}$  would be problematic if any of the gridpoint weights were zero, since the result would then be dividing by  $\mathbf{W}^{1/2}$ . The use of Eq. (4.5) is more robust, and can be used if some of the gridpoint weights are zero. The log pressure-weighted EOF (Fig. 4a) captures the QBO, while the pressure-weighted EOF is dominated by a north-south dipole. The angular momentum EOF (Fig. 4c) captures both features. Although the domain shown in Fig. 4 is that used for the EOF calculation, Eq. (4.1) could be used to extend the plots in height or latitude, or to project a different dataset onto the EOF time series.

## 6. Summary

We have developed a generalized technique for including spatial weighting in EOF-type analyses. EOFs, calculated on a set of (possibly irregular) spatial grid points, should not depend on the grid resolution of spacing. To achieve this goal, we introduce a generalized spatial weighting metric into EOF analysis and show how to then find the resulting generalized EOFs.

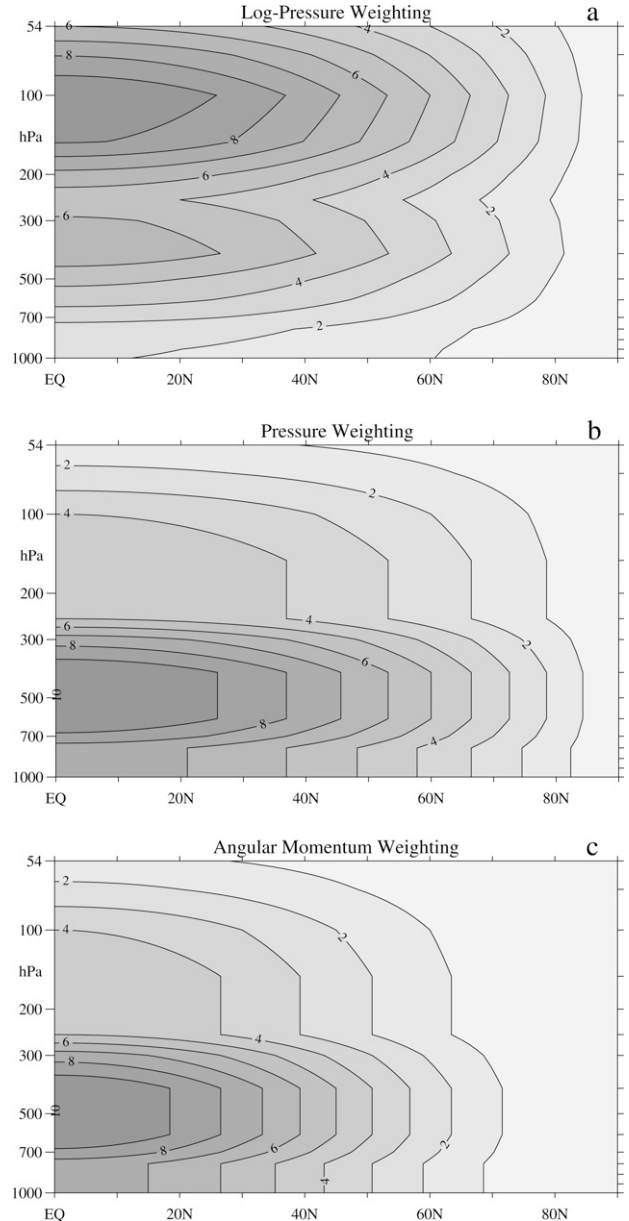


FIG. 3. Weighting functions corresponding to (a) vertical log-(pressure) weighting and (b) vertical pressure weighting. The latitudinal scaling is by  $\cos\theta$  in (a) and (b). (c) Angular momentum weighting with latitudinal scaling by  $\cos^2\theta$ . In all panels, the values have been scaled by an arbitrary scaling factor for plotting. The ticks on the right side of the diagrams show the data levels.

If  $p$  is the spatial dimension of the data matrix, the weighting matrix  $\mathbf{W}$  is dimensioned  $p \times p$ . Its application is simple and straightforward, involving either pretransforming the data by multiplying by  $\mathbf{W}^{1/2}$  or carrying  $\mathbf{W}$  through the EOF calculation.

Although the weighting matrix is developed in order to compensate for an irregular grid, its potential application is much broader. There are a variety of other



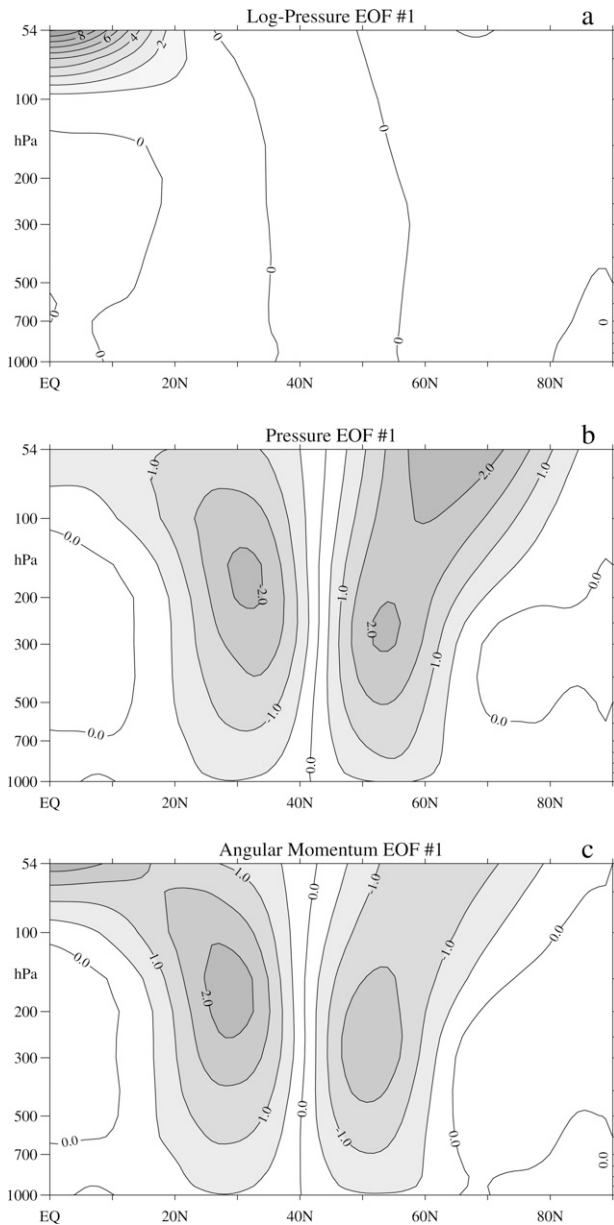


FIG. 4. Leading EOFs of zonal-mean zonal wind, calculated with the weighting functions in Fig. 3: (a) log-pressure weighting, (b) pressure weighting, and (c) angular momentum weighting.

reasons to spatially weight EOF analysis, including the neglect of unrepresented subgrid-scale variance, masking or emphasizing certain regions, compensating for error covariances, etc. In some cases, the weighting matrix can play the role of a filter. For any of these weighting/filtering strategies, the application to EOF analysis is the same.

The weighting techniques are applicable to all variants of EOF analysis such as extended EOFs and maxi-

mum covariance analysis. All these techniques partition the total variance of the data matrix into discrete modes.

The weighting techniques are developed for gridded data that are equally spaced in time, but the same methodology could be applied to the time domain if the observation times are irregular. The results would be nearly identical to interpolating to a grid that is evenly spaced in time.

Fast, memory-efficient methods are required in order to extract EOFs from large datasets. Often, only the first few EOFs are needed, and with a large dataset it can be computationally more efficient to directly calculate the first few EOFs rather than use a standard method that finds all the EOFs at once. This study describes one such approach based on a simple iteration of successive projections of the data onto time series and spatial maps. The method is initialized with a random first guess of the EOF time series and consists of alternately projecting the data matrix onto successive estimates of the EOF and its time series. The technique is a variant of the power method, which has well-defined convergence properties.

The iterative technique works directly from the data matrix (not the covariance matrix) and is suitable for very large data matrices because the covariance matrix is never directly calculated. The technique projects the data onto successive estimates of the EOF spatial pattern and EOF time series. A weighting matrix can be carried through the calculation so that the data are not pre-transformed by multiplying by the square root of appropriate weighting matrices. A subroutine in Interactive Data Language (IDL) that implements this method is available from the authors or online ([www.nwra.com/baldwin](http://www.nwra.com/baldwin)).

We argue that the original spatial resolution of the data should be retained when performing any analysis of patterns in geophysical data. Computational issues may lead researchers to degrade the resolution of the data (e.g., by retaining only every other grid point) to make the calculation tractable (Wallace et al. 1992; class notes on SVD and EOF analyses are available online at <http://jisao.washington.edu/wallace/>). The degraded resolution reduces the total variance of the field. Since the total variance of the field is reduced, the percent of the variance associated with each EOF can change, resulting in patterns that differ from those that would have been calculated from the full data grid. This could also change the order of the EOFs. Degrading the resolution could therefore affect the determination of whether or not empirically determined EOFs are separate (North et al. 1982; Quadrelli et al. 2005).

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## REFERENCES

- Allen, M. R., and L. A. Smith, 1997: Optimal filtering in singular spectrum analysis. *Phys. Lett.*, **234**, 419–428.
- Ambaum, M. H. P., B. J. Hoskins, and D. B. Stephenson, 2001: Arctic Oscillation or North Atlantic Oscillation? *J. Climate*, **14**, 3495–3507.
- Baldwin, M. P., and K.-K. Tung, 1994: Extra-tropical QBO signals in angular momentum and wave forcing. *Geophys. Res. Lett.*, **21**, 2717–2720.
- , and Coauthors, 2001: The quasi-biennial oscillation. *Rev. Geophys.*, **39**, 179–229.
- , D. B. Stephenson, D. W. J. Thompson, T. J. Dunkerton, A. J. Charlton, and A. O'Neill, 2003: Stratospheric memory and skill of extended-range weather forecasts. *Science*, **301**, 636–640.
- Bretherton, C. S., 2003: Empirical orthogonal functions and singular vectors. *Encyclopedia of Atmospheric Sciences*, J. R. Holton, J. A. Pyle, and J. Curry, Eds., Academic Press, 617–620.
- Brunet, G., 1994: Empirical normal mode analysis of atmospheric data. *J. Atmos. Sci.*, **51**, 932–952.
- , and R. Vautard, 1996: Empirical normal modes versus empirical orthogonal functions for statistical prediction. *J. Atmos. Sci.*, **53**, 3468–3489.
- Chang, P., R. Saravanan, L. Ji, and G. C. Hegerl, 2000: The effect of local sea surface temperature on atmospheric circulation over the tropical Atlantic sector. *J. Climate*, **13**, 2195–2216.
- Cheng, X., and J. M. Wallace, 1993: Cluster analysis of the Northern Hemisphere wintertime 500-hPa height field: Spatial patterns. *J. Atmos. Sci.*, **50**, 2674–2696.
- Clint, A., and M. Jennings, 1970: The evaluation of eigenvalues and eigenvectors of real symmetric matrices by simultaneous iteration. *Comput. J.*, **13**, 76–80.
- Craw, I., and P. Cameron, 1992: Face recognition by computer. *Proceedings of the British Machine Vision Conference*, D. Hogg and R. Boyle, Eds., Springer-Verlag, 489–507.
- Folland, C., 1988: The weighting of data in an EOF analysis. MetO13 Discussion Note No. 113, Met Office, Exeter, United Kingdom, 4 pp.
- Golub, G. H., and C. F. Van Loan, 1983: *Matrix Computations*. Johns Hopkins University Press, 476 pp.
- Hannachi, A., I. T. Jolliffe, D. B. Stephenson, and N. Trendafilov, 2006: In search of simple structures in climate: Simplifying EOFs. *Int. J. Climatol.*, **26**, 7–28.
- , —, and —, 2007: Empirical orthogonal functions and related techniques in atmospheric science, a review. *Int. J. Climatol.*, **27**, 1119–1152, doi:10.1002/joc.1499.
- Hawkins, E., and R. Sutton, 2007: Variability of the Atlantic thermal circulation described by three-dimensional empirical orthogonal functions. *Climate Dyn.*, **29**, 745–762, doi:10.1007/s00382-007-0263-8.
- Hewitson, B. C., and R. G. Crane, 2002: Self-organized maps: Applications to synoptic climatology. *Climate Res.*, **22**, 13–26.
- Higham, N. J., 2008: *Functions of Matrices: Theory and Computation*. Society for Industrial and Applied Mathematics, 425 pp.
- Holmström, I., 1963: On a method for parametric representation of the state of the atmosphere. *Tellus*, **15**, 127–149.
- Jolliffe, I. T., 2002: *Principal Component Analysis*. 2nd ed. Springer, 487 pp.
- , 2003: A cautionary note on artificial examples of EOFs. *J. Climate*, **16**, 1084–1086.
- Jongman, R. H. G., C. J. F. ter Braak, and O. F. R. van Tongeren, Eds., 1995: *Data Analysis in Community and Landscape Ecology*. Cambridge University Press, 299 pp.
- Kagan, R. L., L. S. Gandin, and T. M. Smith, 1997: *Averaging of Meteorological Fields*. Kluwer Academic Publishers, 279 pp.
- Legendre, P., and L. Legendre, 1998: *Numerical Ecology*. 2nd ed. Elsevier, 853 pp.
- Monin, A. S., and A. M. Yaglom, 1971: *Statistical Fluid Mechanics*. Vol. 1, *Mechanics of Turbulence*, MIT Press, 769 pp.
- North, G. R., T. L. Bell, R. F. Cahalan, and F. J. Moeng, 1982: Sampling errors in the estimation of empirical orthogonal functions. *Mon. Wea. Rev.*, **110**, 699–706.
- Quadrelli, R., C. S. Bretherton, and J. M. Wallace, 2005: On sampling errors in empirical orthogonal functions. *J. Climate*, **18**, 3704–3710.
- Reusch, D. B., R. B. Alley, and B. C. Hewitson, 2007: North Atlantic climate variability from a self-organizing map perspective. *J. Geophys. Res.*, **112**, D02104, doi:10.1029/2006JD007460.
- Stephenson, D. B., 1997: Correlation of spatial climate/weather maps and the advantages of using the Mahalanobis metric in predictions. *Tellus*, **49A**, 513–527.
- Strang, G., 2003: *Introduction to Linear Algebra*. 3rd ed. Wellesley Cambridge Press, 568 pp.
- Thompson, D. W. J., and J. M. Wallace, 1998: The Arctic Oscillation signature in the wintertime geopotential height and temperature fields. *Geophys. Res. Lett.*, **25**, 1297–1300.
- Uppala, S. M., and Coauthors, 2005: The ERA-40 Re-Analysis. *Quart. J. Roy. Meteor. Soc.*, **131**, 2961–3012.
- Van den Dool, H. M., 2007: *Empirical Methods in Short-Term Climate Prediction*. Oxford University Press, 215 pp.
- , S. Saha, and Å. Johansson, 2000: Empirical orthogonal teleconnections. *J. Climate*, **13**, 1421–1435.
- Venzke, S., M. R. Allen, R. T. Sutton, and D. P. Rowell, 1999: The atmospheric response over the North Atlantic to decadal changes in sea surface temperature. *J. Climate*, **12**, 2562–2584.
- von Storch, H., and F. W. Zwiers, 1999: *Statistical Analysis in Climate Research*. Cambridge University Press, 484 pp.
- von Storch, J.-S., 2000: Angular momenta of the Antarctic and the Arctic Oscillations. *J. Climate*, **13**, 681–685.
- Wallace, J. M., C. Smith, and C. S. Bretherton, 1992: Singular value decomposition of wintertime sea surface temperature and 500-mb height anomalies. *J. Climate*, **5**, 561–576.
- , Y. Zhang, and K.-H. Lau, 1993a: Structure and seasonality of interannual and interdecadal variability of the geopotential height and temperature fields in the Northern Hemisphere. *J. Climate*, **6**, 2063–2082.
- , R. L. Panetta, and J. Estberg, 1993b: Representation of the equatorial stratospheric quasi-biennial oscillation in EOF phase space. *J. Atmos. Sci.*, **50**, 1751–1762.